

Asymptotics of certain families of Higgs bundles

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(joint with Brian Collier, UIUC)

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Composing with the the unique irreducible representation $PSL(2, \mathbb{R}) \rightarrow PSL(n, \mathbb{R})$ we obtain a distinguished component of $Rep_{PSL(n, \mathbb{R})}$, called *Hitchin component*.

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The proof of this theorem uses Higgs bundle techniques. The Higgs bundle parametrization of $Hit_n(S)$ is as follows:

$$\text{Bundle: } E = K^{\frac{n-1}{2}} \oplus K^{\frac{n-3}{2}} \oplus \dots \oplus K^{-\frac{n-3}{2}} \oplus K^{-\frac{n-1}{2}}$$

$$\text{Higgs field: } \phi = \begin{pmatrix} 0 & \frac{n-1}{2}q_2 & & \dots & q_{n-1} & q_n \\ 1 & 0 & \frac{n-3}{2}q_2 & \dots & q_{n-2} & q_{n-1} \\ & \ddots & & \ddots & & \\ & & & & \frac{n-3}{2}q_2 & \\ & & & 1 & 0 & \frac{n-1}{2}q_2 \\ & & & & 1 & 0 \end{pmatrix}$$

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Conversely, if (A_h, ϕ) solves Hitchin equation, then the Higgs bundle (E, ϕ) is polystable.

If (A_h, ϕ) solves Hitchin equation, then $A_h + \phi + \phi^{*h}$ is a **flat** $SL(n, \mathbb{C})$ -connection.

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- To answer the above questions, the first step is to solve Hitchin equation asymptotically. This is in general impossible.
- However, we manage to understand the two cases:
 - (1) $t(0, \dots, q_n)$
 - (2) $t(0, \dots, q_{n-1}, 0)$.

Why these two cases?

Theorem (Baraglia n , Collier $n - 1$)

For $(0, \dots, q_n)$ and $(0, \dots, q_{n-1}, 0)$, the metric solving Hitchin equation is diagonal on the line bundles

$$K^{\frac{n-1}{2}} \oplus K^{\frac{n-3}{2}} \oplus \dots \oplus K^{-\frac{n-3}{2}} \oplus K^{-\frac{n-1}{2}}$$

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$$\begin{cases} F_{A_1} + t^2 h_1^2 q_n \wedge \bar{q}_n - h_1^{-1} h_2 = 0 \\ F_{A_j} + h_{j-1}^{-1} h_j - h_j^{-1} h_{j+1} = 0 & 1 < j < \frac{n}{2} \\ F_{A_{\frac{n}{2}}} + h_{\frac{n}{2}-1}^{-1} h_{\frac{n}{2}} - h_{\frac{n}{2}}^{-2} = 0 \end{cases}$$

Theorem (Collier-L)

For $(0, \dots, tq_n) \in \bigoplus_{j=2}^n H^0(K^j)$, at any point $p \in \Sigma$ away from the zeros of q_n , as $t \rightarrow \infty$, the metric $h_j(t)$ on $K^{\frac{n+1-2j}{2}}$ admits the expansion

$$h_j(t) = (t|q_n|)^{-\frac{n+1-2j}{n}} \left(1 + O\left(t^{-\frac{2}{n}}\right) \right) \quad \text{for all } j$$

The analogous result is also true for $(0, \dots, tq_{n-1}, 0)$.

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- Note that the Hitchin equation is highly nontrivial. The solutions are **globally** depending on the parameters q_n . Here, our results show that asymptotically, the solutions to Hitchin system only depend on the **local** values of q_n .

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- Main Tool: the maximum principle (numerous times)

Parallel Transport Asymptotics

Given a path $\gamma : [0, L] \rightarrow \tilde{\Sigma}$, let $T_\gamma(t) : E_{\gamma(0)} \rightarrow E_{\gamma(L)}$ be parallel transport operators along γ for flat connections ∇_t .

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$$T_\gamma(t) = \left(Id + O\left(t^{-\frac{1}{2n}}\right) \right) S \begin{pmatrix} e^{-t^{\frac{1}{n}} \mu_1} & & \\ & \ddots & \\ & & e^{-t^{\frac{1}{n}} \mu_n} \end{pmatrix} S^{-1}$$

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where S is constant unitary, and $\mu_k = 2\operatorname{Re} \left(\int_\gamma e^{\frac{2\pi ki}{n}} z \right)$.

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- This theorem relies on very technical analysis on the error estimates of the Hermitian metric solution and the special structure of Toda lattice.
- Both theorems are proved in $(0, q_3)$ case by J. Loftin.

Work in Progress (Quiver bundles)

Instead of thinking Higgs bundle

$$\mathcal{E} = K^{\frac{n-1}{2}} \oplus K^{\frac{n-3}{2}} \oplus \dots \oplus K^{-\frac{n-3}{2}} \oplus K^{-\frac{n-1}{2}}$$

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Application: Understanding asymptotics of maximal representation in $Sp(4, \mathbb{R})$.

Thank You!