

Chaos in periodically forced reversible vector fields

Isabel S. Labouriau and Elisa Sovrano

Abstract. We discuss the appearance of chaos in time-periodic perturbations of reversible vector fields in the plane. We use the normal forms of codimension 1 reversible vector fields and discuss the ways a time-dependent periodic forcing term of pulse form may be added to them to yield topological chaotic behaviour. Chaos here means that the resulting dynamics is semi-conjugate to a shift in a finite alphabet.

The results rely on the classification of reversible vector fields and on the theory of topological horseshoes. This work is part of a project of studying periodic forcing of symmetric vector fields.

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1. Introduction

A standard classification of continuous dynamical systems defined by a set of first order ordinary differential equations distinguishes between conservative systems and dissipative ones [10]. On the one hand, conservative systems can be described by a Hamiltonian function. By varying the initial conditions, these systems can exhibit regions of regular motions surrounded by a sea of chaotic ones. Instead, dealing with dissipative systems, conserved quantities are no longer guaranteed, and chaotic regions could coexist with stable equilibria, limit cycles, and strange attractors.

In between conservative and dissipative systems, there are systems with reversing symmetries. By reversible dynamical systems we mean those admitting an involution in phase space which reverses the direction of time (see [1, 3, 8, 11]). It is shown that these systems despite having similar features to Hamiltonian ones (e.g., at an elliptic equilibrium can possess the same structure), yet they are different because they can also have attractors and repellers.

The additional structure given by reversing symmetries allows exhibiting complex behaviors for codimension one bifurcations, and so, it can be responsible for chaotic dynamics.

The goal of this paper is to find chaos for a class of planar periodically perturbed reversible systems whose normal form analysis is studied in [11]. We take into account the local bifurcations of low codimension by arguing what dynamical behaviors we can expect. Our main result is the following.

Conjecture 1.1. *Let $X_\lambda(x, y)$ be a fixed type of normal form for a one-parameter family of codimension 1 reversible vector fields. Let λ_1 and λ_2 be two real distinct values. Suppose that the dynamical system $\dot{X} = X(x, y)$ switches in a T -periodic manner between*

$$\dot{X} = X_{\lambda_1}(x, y) \text{ for } t \in [0, \tau_1) \quad \text{and} \quad \dot{X} = X_{\lambda_2}(x, y) \text{ for } t \in [\tau_1, \tau_2)$$

with $\tau_1 + \tau_2 = T$. Then for open sets of the parameters (λ_1, λ_2) and for τ_1 and τ_2 in open intervals there exist infinitely many T -periodic solutions as well as chaotic-like dynamics for the problem $\dot{X} = X(x, y)$.

The paper is organized as follows. In Section 2 we discuss the classification of plane reversible vector fields of codimension 0 and 1. In Section 3 we give a review of the concept of symbolic dynamics and topological horseshoes. We collect preliminary topological results in the phase-plane that can produce chaotic dynamics. In Section 4 we prove Conjecture 1.1 for the two of the four normal forms of codimension 1 reversible vector fields: *i*) saddle type and *ii*) cusp type. We conjecture that the other two possible normal forms, namely *iii*) nodal type and *iv*) focal type, may also be amenable to the same treatment.

2. Planar reversible systems

In [11], M. A. Teixeira has provided a local classification of 2D reversible systems of codimension less than or equal to two. A dynamical system $\dot{X} = V(X)$ is called *reversible* if there is a phase space *involution* h (i.e., $h^2 = \text{Id}$) such that $Dh(p)V(p) = -V(h(p))$ for $p \in \mathbb{R}^2$. We deal with reversible planar systems where the involution is $h(x, y) = (x, -y)$. Hence, we consider a dynamical system of the following form

$$\begin{cases} \dot{x} = yf(x, y^2), \\ \dot{y} = g(x, y^2), \end{cases} \quad (2.1)$$

where the functions f and g are smooth. We consider the behaviour of (2.1) near the origin, often making the assumption that it has an equilibrium at the origin. In the half-plane $y > 0$, by using the transformation $u = x$ and $v = y^2$, we can write system (2.1) equivalently as follows

$$\begin{cases} \dot{u} = \sqrt{v}f(u, v), \\ \dot{v} = 2\sqrt{v}g(u, v). \end{cases}$$

Through the symmetry properties of the vector field $X(x, y)$ associated with (2.1), the behavior of X near $(0, 0)$ may be described by the analysis in the half-plane $\{(u, v) \in \mathbb{R}^2 : v \geq 0\}$ of the vector field $Y(u, v) = (f(u, v), g(u, v))$.

2.1. Normal forms

Following the work in [11], the generic equilibria of reversible ODEs near the origin are either centers and saddles on the line of symmetry or a couple of repellers and attractors, as in Figure 1.

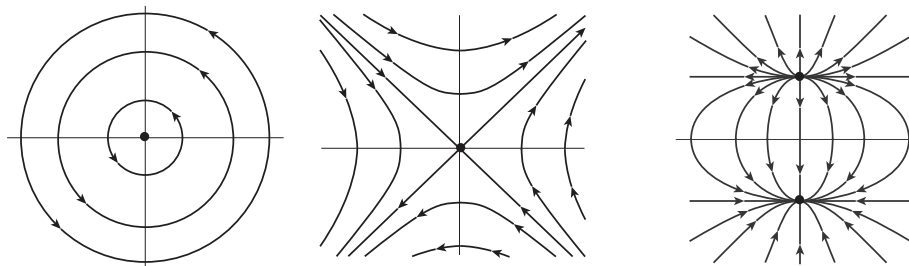


FIGURE 1. Phase-portraits of equilibria occurring in generic 2D reversible fields. The local geometry may be of a center (left), a saddle (middle), or a pair of attractor and repeller (right).

Let S be the line $\{(x, 0) : x \in \mathbb{R}\}$, the set of fixed points for h . An equilibrium point of V that lies on S is called a *symmetric equilibrium*.

Theorem 2.1 ([11]). *The normal forms around a symmetric equilibrium at $(0, 0)$ of a structurally stable reversible vector field X are:*

- $X(x, y) = (y, x)$,
- $X(x, y) = (y, -x)$.

In the first case the origin is a center, and in the second one it is a saddle. The next result classifies one parameter families X_λ of reversible vector fields such that X_0 has a symmetric equilibrium at the origin.

Theorem 2.2 ([11]). *The normal forms of one-parameter families of structurally stable reversible vector fields X_λ near a symmetric equilibrium at $(0, 0)$ are:*

- i) *saddle type:* $X_\lambda(x, y) = (xy, x - y^2 + \lambda)$,
- ii) *cusp type:* $X_\lambda(x, y) = (y, x^2 + \lambda)$,
- iii) *nodal type:* $X_\lambda(x, y) = (xy, x + 2y^2 + \lambda)$ or $X_\lambda(x, y) = (-xy, x - 2y^2 + \lambda)$,
- iv) *focal type:* $X_\lambda(x, y) = (xy + y^3, -x + y^2 + \lambda)$.

Depending on λ , the phase-portraits of the above normal forms can be described as follows.

Figure 2 shows the phase portraits of the *saddle type*. When $\lambda \leq 0$ there is an equilibrium at $(-\lambda, 0)$ which is a saddle. When $\lambda > 0$ there are three

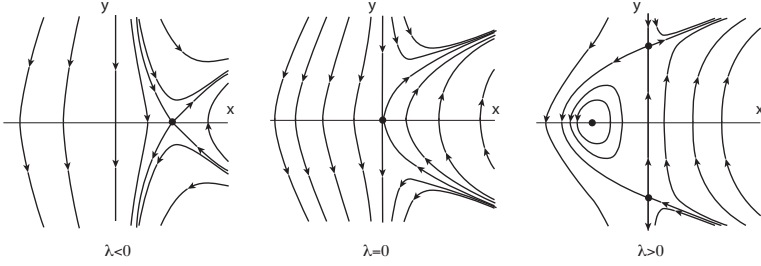


FIGURE 2. Phase-portraits reversible vector fields of saddle type.

equilibria: a center and two saddles at $(-\lambda, 0)$, $(0, -\sqrt{\lambda})$ and $(0, \sqrt{\lambda})$, respectively. The saddle points are connected through heteroclinic trajectories which surround periodic orbits.

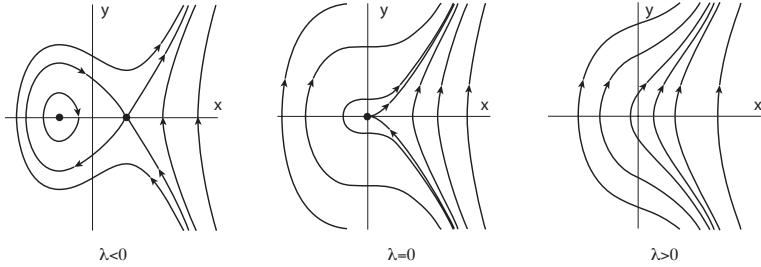


FIGURE 3. Phase-portraits reversible vector fields of cusp type.

Concerning the *cusp type* when $\lambda < 0$ there are two equilibria: a center and a saddle which are at $(-\sqrt{-\lambda}, 0)$ and $(\sqrt{-\lambda}, 0)$, respectively. Due to the reversibility, the only periodic orbits are the ones that meet the points $(x, 0)$ with $-2\sqrt{-\lambda} < x < \sqrt{-\lambda}$, as in Figure 3. Moreover, these orbits are located inside the homoclinic trajectory that passes through $(-2\sqrt{-\lambda}, 0)$. When $\lambda = 0$ there is only an equilibrium which is a degenerate saddle at $(0, 0)$ and all the orbits are unbounded. When $\lambda > 0$ there are no equilibria.

For the *nodal type* (first case, shown in Figure 4) when $\lambda < 0$ there are three equilibria: an attractor, a repeller and a saddle at $(-\sqrt{-\lambda/2}, 0)$, $(\sqrt{-\lambda/2}, 0)$ and $(-\lambda, 0)$, respectively. When $\lambda = 0$ there is only an equilibrium at $(0, 0)$. When $\lambda > 0$ there is only an equilibrium at $(-\lambda, 0)$ which is a center and all the orbits are periodic. In the second case there is always an equilibrium at $(-\lambda, 0)$ and for $\lambda > 0$ there is also a pair of equilibria at $(0, \pm\sqrt{\lambda})$.

For the *focal type* when $\lambda < 0$ there are three equilibria: a saddle and two foci at $(\lambda, 0)$, $(\lambda/2, -\sqrt{\lambda/2})$ and $(\lambda/2, \sqrt{\lambda/2})$, respectively. When $\lambda \geq 0$ there is only an equilibrium at $(-\lambda, 0)$ which is a center and all the orbits are periodic as in Figure 5.

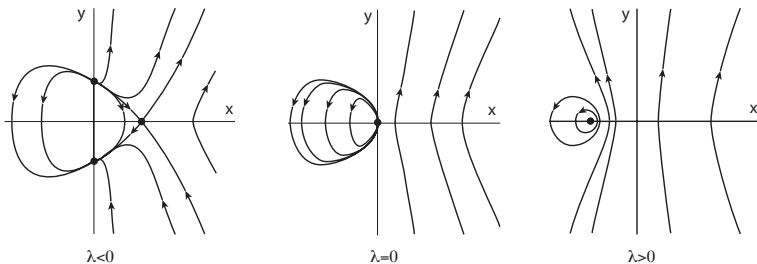


FIGURE 4. Phase-portraits reversible vector fields of nodal type (first case).

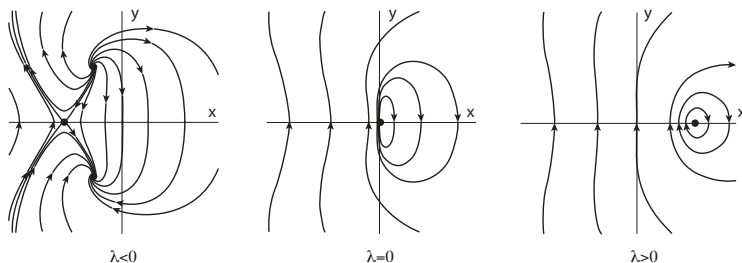


FIGURE 5. Phase-portraits reversible vector fields of focal type.

3. Background on chaotic dynamics and preliminary results

3.1. Symbolic dynamics and chaos

To review the topological approach exploited throughout the paper, we start by introducing some notation and definitions of symbolic dynamics.

Let $\Sigma_m := \{0, \dots, m - 1\}^{\mathbb{Z}}$ be the set of all two-sided sequences $S = (s_i)_{i \in \mathbb{Z}}$ with $s_i \in \{0, \dots, m - 1\}$ for each $i \in \mathbb{Z}$ endowed with a standard metric that makes Σ_m a compact space with the product topology. We define the shift map $\sigma : \Sigma_m \rightarrow \Sigma_m$ by $\sigma(S) = S' = (s'_i)_{i \in \mathbb{Z}}$ with $s'_i = s_{i+1}$ for all $i \in \mathbb{Z}$. We say that a map h on a metric space is semiconjugate (respectively, conjugate) to the shift map on m symbols if there exists a compact invariant set Λ and a continuous and surjective (respectively, bijective) map $\Pi : \Lambda \rightarrow \Sigma_m$ such that $\Pi \circ h(w) = \sigma \circ \Pi(w)$, for all $w \in \Lambda$.

The deterministic chaos is usually associated with the possibility to reproduce all the possible outcomes of a coin-tossing experiment, by varying the initial conditions within the dynamical system. We can express this concept using the symbolic dynamics of the shift map on the sets of two-sided sequences of 2 symbols. However, by considering a finite alphabet made by m symbols the possible dynamics can be more complex. Hence, in the sequel we adopt the following definition of chaos (cf., [4, 5]).

Definition 3.1 (Symbolic dynamics). *Let $h : \text{dom } h \subseteq \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be a map and let $\mathcal{D} \subseteq \text{dom } h$ be a nonempty set. We say that h induces chaotic dynamics on*

$m \geq 2$ symbols on a set \mathcal{D} if there exist m nonempty pairwise disjoint compact sets $\mathcal{K}_0, \dots, \mathcal{K}_{m-1} \subseteq \mathcal{D}$ such that for each two-sided sequence $(s_i)_{i \in \mathbb{Z}} \in \Sigma_m$ there exists a corresponding sequence $(w_i)_{i \in \mathbb{Z}} \in \mathcal{D}^{\mathbb{Z}}$ such that

$$w_i \in \mathcal{K}_{s_i} \text{ and } w_{i+1} = h(w_i) \text{ for all } i \in \mathbb{Z}, \quad (3.1)$$

and, whenever $(s_i)_{i \in \mathbb{Z}}$ is a k -periodic sequence for some $k \geq 1$ there exists a k -periodic sequence $(w_i)_{i \in \mathbb{Z}} \in \mathcal{D}^{\mathbb{Z}}$ satisfying (3.1).

For a one-to-one map h , Definition 3.1 ensures the existence of a nonempty compact invariant set $\Lambda \subseteq \bigcup_{i=0}^{m-1} \mathcal{K}_i \subseteq \mathcal{D}$ and a continuous surjection Π such that $h|_{\Lambda}$ is semiconjugate to the Bernoulli shift map on $m \geq 2$ symbols. Moreover, it guarantees that the set of the periodic points of h is dense in Λ and, for all two-sided periodic sequences $S \in \Sigma_m$, the preimage $\Pi^{-1}(S)$ contains a periodic point of h with the same period. In this respect Definition 3.1 is related to the concept of topological horseshoe introduced in [2]. This is a weaker notion of chaos than the Smale's horseshoe because the latter requires the full conjugacy between $h|_{\Lambda}$ and the shift map on m symbols.

We introduce the notion of an oriented topological rectangle and the stretching along the path property by borrowing the notations and definitions from [4, 6]. The pair $\widehat{\mathcal{R}} := (\mathcal{R}, \mathcal{R}^-)$ is called *oriented topological rectangle* if $\mathcal{R} \subseteq \mathbb{R}^2$ is a set homeomorphic to $[0, 1] \times [0, 1]$, and $\mathcal{R}^- = \mathcal{R}_l^- \cup \mathcal{R}_r^-$, where \mathcal{R}_l^- and \mathcal{R}_r^- are two disjoint compact arcs contained in $\partial \mathcal{R}$.

Definition 3.2 (SAP property). *Given two topological oriented rectangles $\widehat{\mathcal{R}}_1 := (\mathcal{R}_1, \mathcal{R}_1^-)$, $\widehat{\mathcal{R}}_2 := (\mathcal{R}_2, \mathcal{R}_2^-)$ and a continuous map $h : \text{dom } h \subseteq \mathbb{R}^2 \rightarrow \mathbb{R}^2$, we say that h stretches $\widehat{\mathcal{R}}_1$ to $\widehat{\mathcal{R}}_2$ along the paths if there exists a compact subset \mathcal{K} of $\mathcal{R}_1 \cap \text{dom } h$ and for each path $\gamma : [0, 1] \rightarrow \mathcal{R}_1$ such that $\gamma(0) \in \mathcal{R}_{1,l}^-$ and $\gamma(1) \in \mathcal{R}_{1,r}^-$ (or vice-versa), there exists $[t_0, t_1] \subseteq [0, 1]$ such that*

- $\gamma(t) \in \mathcal{K}$ for all $t \in [t_0, t_1]$,
- $h(\gamma(t)) \in \mathcal{R}_2$ for all $t \in [t_0, t_1]$,
- $h(\gamma(t_0))$ and $h(\gamma(t_1))$ belong to different components of \mathcal{R}_2^- .

In this case, we write

$$(\mathcal{K}, h) : \widehat{\mathcal{R}}_1 \rightleftarrows \widehat{\mathcal{R}}_2.$$

Given a positive integer m , we say that h stretches $\widehat{\mathcal{R}}_1$ to $\widehat{\mathcal{R}}_2$ along the paths with crossing number m and we write

$$h : \widehat{\mathcal{R}}_1 \rightleftarrows^m \widehat{\mathcal{R}}_2$$

if there exist m pairwise disjoint compact sets $\mathcal{K}_0, \dots, \mathcal{K}_{m-1} \subseteq \mathcal{R}_1 \cap \text{dom } h$ such that $(\mathcal{K}_i, h) : \widehat{\mathcal{R}}_1 \rightleftarrows \widehat{\mathcal{R}}_2$ for each $i \in \{0, \dots, m-1\}$.

Finally, in order to detect chaos, a useful topological tool is the Stretching Along the Paths (SAP) method introduced in [5]. In our framework, it can be stated as follows (cf., [4, Th. 2.1]).

Theorem 3.1 (SAP method). *Let $h_1 : \text{dom } \nu \subseteq \mathbb{R}^2 \rightarrow \mathbb{R}^2$ and $h_2 : \text{dom } \eta \subseteq \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be continuous maps. Let $\widehat{\mathcal{R}}_1 = (\mathcal{R}_1, \mathcal{R}_1^-)$ and $\widehat{\mathcal{R}}_2 = (\mathcal{R}_2, \mathcal{R}_2^-)$ be two oriented rectangles in \mathbb{R}^2 . Suppose that*

- there exist $n \geq 1$ pairwise disjoint compact subsets of $\mathcal{R}_1 \cap \text{dom } \nu$, $\mathcal{Q}_0, \dots, \mathcal{Q}_{n-1}$, such that $(\mathcal{Q}_i, h_1): \widehat{\mathcal{R}}_1 \xrightarrow{\cong} \widehat{\mathcal{R}}_2$ for $i = 0, \dots, n-1$,
- there exist $m \geq 1$ pairwise disjoint compact subsets of $\mathcal{R}_2 \cap \text{dom } \eta$, $\mathcal{K}_0, \dots, \mathcal{K}_{m-1}$, such that $(\mathcal{K}_i, h_2): \widehat{\mathcal{R}}_2 \xrightarrow{\cong} \widehat{\mathcal{R}}_1$ for $i = 0, \dots, m-1$.

If at least one between n and m is greater or equal than 2, then the map $h = h_2 \circ h_1$ induces chaotic dynamics on $n \times m$ symbols on

$$\mathcal{Q}^* = \bigcup_{\substack{i=1, \dots, n \\ j=1, \dots, m}} \mathcal{Q}_i \cap \nu^{-1}(\mathcal{K}_j).$$

For the proof of Theorem 3.1 we refer to [4, Th. 3.2].

3.2. Topological tools in the phase-plane

The geometry associated to the phase-portrait of (2.1) exhibits unbounded solutions and periodic trajectories. These configurations guarantee the existence of two types of invariant regions: topological strips and topological annuli confined between unbounded and bounded solutions, respectively. In this section we will give some preliminary topological results on the phase-plane (x, y) needed to establish the dynamics induced by (2.1).

By a *topological strip* \mathcal{S} we mean the image of a straight strip of finite width $\mathbf{S} := \{(x, y) \in \mathbb{R}^2 : x \in \mathbb{R}, -1 \leq y \leq 1\}$ through a local homeomorphism

$$h_{\mathbf{S}}: (0, 1) \times [-1, 1] \rightarrow \mathcal{S}.$$

Let a *bridge* in \mathcal{S} be the image by $h_{\mathbf{S}}$ of any simple continuous curve $\gamma: [a, b] \rightarrow \mathbf{S}$ such that $\gamma(a) = -1$ and $\gamma(b) = 1$ or, viceversa, $\gamma(a) = 1$ and $\gamma(b) = -1$.

A *topological annulus* \mathcal{A} is defined as the image of a rectangular region $\mathbf{A} := \{(x, y) \in \mathbb{R}^2 : 1 \leq x \leq 2, -1 \leq y \leq 1\}$ through a continuous map

$$h_{\mathbf{A}}: [1, 2] \times [-1, 1] \rightarrow \mathcal{A},$$

such that the restriction of $h_{\mathbf{A}}$ to $(1, 2) \times [-1, 1]$ is a homeomorphism and $h_{\mathbf{A}}(1, y) = h_{\mathbf{A}}(2, y)$. We notice that the restriction to $(1, 2) \times [-1, 1]$ yields a strip. Moreover, the boundary of the topological annulus $\partial\mathcal{A}$ is the union of two Jordan curves $\partial^i\mathcal{A} := h_{\mathbf{A}}(x, -1)$ and $\partial^e\mathcal{A} := h_{\mathbf{A}}(x, 1)$. We denote the portion of the plane outside a generic Jordan curve Γ by $out(\Gamma)$ and the one inside by $in(\Gamma)$. For identification purposes, let $\partial^i\mathcal{A} \subset in(\partial^e\mathcal{A})$. In this manner, we can identify two connected sets, one bounded and another one unbounded given by $in(\partial^i\mathcal{A})$ and $out(\partial^e\mathcal{A})$, respectively. Let a *ray* in \mathcal{A} be any simple continuous curve $\gamma: [a, b] \rightarrow \mathcal{A}$ such that $\gamma(a) \in \partial^i\mathcal{A}$ and $\gamma(b) \in \partial^e\mathcal{A}$ or, viceversa, $\gamma(a) \in \partial^e\mathcal{A}$ and $\gamma(b) \in \partial^i\mathcal{A}$.

We are interested in crossing configurations between either an annulus and a strip or two annuli. In particular we are looking for similarities with the geometry of the linked-twist maps (see [7, 12]). Hence, we introduce the following definition and in Figure 6 we provide a visual representation of the linkage condition between an annulus and a strip.

Definition 3.3 (Linkage condition). *Let \mathcal{A} be a topological annulus and \mathcal{S} be a topological strip. We say that \mathcal{A} is linked with \mathcal{S} if there exist a bridge γ_1 in \mathcal{S} , a ray γ_2 in \mathcal{A} , and a topological ball B containing \mathcal{A} such that:*

- $\gamma_1 \subset \text{in}(\partial^i \mathcal{A})$;
- $\gamma_2 \cap \mathcal{S} = \emptyset$;
- $(\mathcal{S} \setminus \gamma_1) \cap \partial B$ consists of exactly two disjoint bridges.

From Definition 3.3 we observe that when \mathcal{A} is linked with \mathcal{S} , then the topological ball B is cut into two connected components B^+ and B^- .

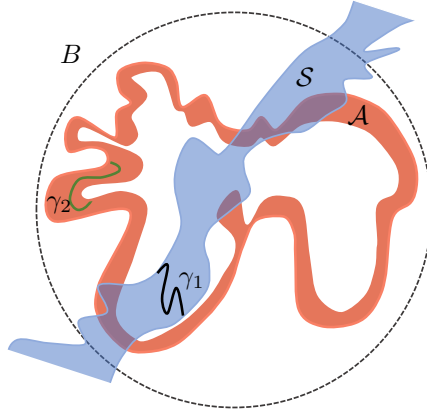


FIGURE 6. Linkage condition. The figure represents an example of a topological annulus (red) linked with a topological strip (blue) through the existence of a bridge (black) and a ray (green).

Notice that Definition 3.3 involves only the geometry inside a topological ball B . Therefore it could include the case when the strip \mathcal{S} is the intersection of an annulus \mathcal{A}_2 with the ball B . In this manner we are generalizing the definition of the linkage between two annuli $\mathcal{A}_1, \mathcal{A}_2$ given in [6]. In the following proposition we also recover some of the properties collected in [6, Proposition 3.1] for the linkage of two annuli.

From the third requirement of Definition 3.3 it follows that the set $B \setminus \mathcal{S}$ has two connected components that will be denoted B^+ and B^- .

Proposition 3.2. *If the topological strip \mathcal{S} is linked with the topological annulus \mathcal{A} , then there exists a topological ball B containing \mathcal{A} , a bridge γ_3 in \mathcal{S} and a ray γ_4 in \mathcal{A} such that $\gamma_3 \subset B \setminus \text{in}(\partial^e \mathcal{A})$, and denoting by B^+ the component of $B \setminus \mathcal{S}$ that contains $\gamma_2 \subset B^+$, then $\gamma_4 \subset B^-$.*

Proof. First of all we observe that the existence of a bridge $\gamma_3 \subset B \setminus \text{in}(\partial^e \mathcal{A})$ follows immediately from Definition 3.3. Indeed, we can choose γ_3 between one of the two components of $(\mathcal{S} \setminus \gamma_1) \cap \partial B$ and one of the bridges in $(\mathcal{S} \setminus \gamma_1) \cap \partial B$.

The proof of the existence of the ray γ_4 is entirely analogous to that of [6, Proposition 3.1] and is omitted. \square

In the sequel, we deal with the study of the dynamics in a strip \mathcal{S} and in an annulus \mathcal{A} . If they are linked, then there exist two disjoint topological rectangular regions $\mathcal{R}_1 \subset \mathcal{A} \cap \mathcal{S} \cap B$ and $\mathcal{R}_2 \subset \mathcal{A} \cap \mathcal{S} \cap B$.

Firstly, we consider the following continuous map

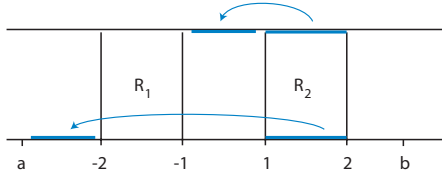
$$\phi_{\mathcal{S}}: \mathcal{S} \rightarrow \mathcal{S}. \tag{3.2}$$

Without loss of generality, we can assume that $\mathcal{R}_1, \mathcal{R}_2$ are homeomorphic to $R_1 = [-2, -1] \times [-1, 1]$ and $R_2 = [1, 2] \times [-1, 1]$, respectively. We suppose that the map $\phi_{\mathcal{S}}$ in (3.2) admits a lift $\tilde{\phi}_{\mathcal{S}}$ to the covering space $[a, b] \times [-1, 1]$, with $a < -2$ and $b > 2$, defined as

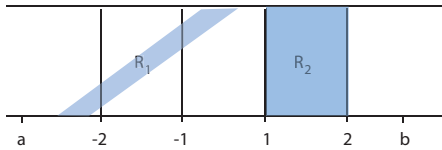
$$\tilde{\phi}_{\mathcal{S}}: (x, y) \mapsto (x + \Xi(x, y), \zeta(x, y)),$$

where ζ, Ξ are continuous functions.

Definition 3.4 (Strip boundary invariance condition). *The condition holds for the map $\phi_{\mathcal{S}}$ if the second coordinate of its lift $\tilde{\phi}_{\mathcal{S}}$ satisfies $\zeta(x, -1) \equiv -1$ and $\zeta(x, 1) \equiv 1$.*



(A) Image of $[1, 2] \times \{-1\}$ and $[1, 2] \times \{1\}$ under a twist condition with respect to the rectangle R_1 .



(B) Image of the rectangle R_2 under a twist condition with respect to the rectangle R_1 .

FIGURE 7. Example of strip twist condition.

Definition 3.5 (Strip twist condition). *The condition holds with respect to R_1 for $x \in [1, 2]$ if either*

$$\Xi(x, -1) \leq -4 \quad \text{and} \quad \Xi(x, 1) \geq -2,$$

or

$$\Xi(x, -1) \geq -2 \quad \text{and} \quad \Xi(x, 1) \leq -4.$$

The condition holds with respect to R_2 for $x \in [-2, -1]$ if either

$$\Xi(x, -1) \leq 2 \quad \text{and} \quad \Xi(x, 1) \geq 4,$$

or

$$\Xi(x, -1) \geq 4 \quad \text{and} \quad \Xi(x, 1) \leq 2.$$

Secondly, we consider the following continuous map

$$\phi_{\mathcal{A}}: \mathcal{A} \rightarrow \mathcal{A}. \quad (3.3)$$

We suppose that the map $\phi_{\mathcal{A}}$ in (3.3) admits a lift $\tilde{\phi}_{\mathcal{A}}$ to the covering space $\mathbb{R} \times [-1, 1]$ defined as

$$\tilde{\phi}_{\mathcal{A}}: (\theta, \rho) \mapsto (\theta + \Theta(\theta, \rho), \omega(\theta, \rho)),$$

where θ, ρ are generalized polar coordinates, and Θ, ω are continuous functions 1-periodic in the θ -variable. Without loss of generality, we can assume that \mathcal{R}_1 and \mathcal{R}_2 are represented in the covering by $R_1 = [2k, 2k + \frac{1}{2}] \times [-1, 1]$ and $R_2 = [2k + 1, 2k + \frac{3}{2}] \times [-1, 1]$, respectively.

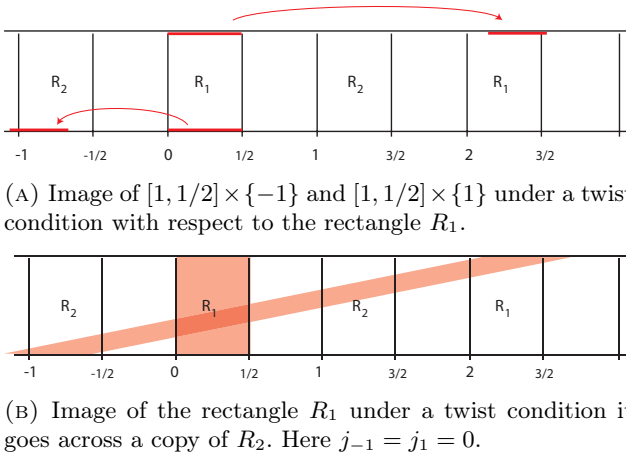


FIGURE 8. Example of an annular twist condition.

Definition 3.6 (Annular boundary invariance condition). *The condition holds for the map $\phi_{\mathcal{A}}$ if the second coordinate of its lift $\tilde{\phi}_{\mathcal{A}}$ satisfies $\omega(\theta, -1) \equiv -1$ and $\omega(\theta, 1) \equiv 1$.*

Definition 3.7 (Annular twist condition). *There exist integers j_{-1} and j_1 such that the condition holds with respect to R_1 for $\theta \in [0, 1/2]$ if either*

$$\Theta(\theta, -1) \leq 2j_{-1} + \frac{1}{2} \quad \text{and} \quad \Theta(\theta, 1) \geq 2j_1 + \frac{3}{2}, \quad \text{with } j_1 + 1 - j_{-1} > 0$$

or

$$\Theta(\theta, -1) \geq 2j_{-1} + \frac{3}{2} \quad \text{and} \quad \Theta(\theta, 1) \leq 2j_1 + \frac{1}{2}, \quad \text{with } j_{-1} + 1 - j_1 > 0$$

hold.

We notice that when the annular twist condition holds with respect to R_1 then the rectangle R_1 is stretched across R_2 a number of times which is given by $|j_{-1} - j_1| + 1$.

Theorem 3.3. *Let \mathcal{A} be a topological annulus linked with a topological strip \mathcal{S} . Let \mathcal{R}_i for $i = 1, 2$ be two disjoint oriented topological rectangles given through the linkage. Let $\phi_{\mathcal{A}}: \mathcal{A} \rightarrow \mathcal{A}$ and $\phi_{\mathcal{S}}: \mathcal{S} \rightarrow \mathcal{S}$, be two continuous maps that satisfy the boundary invariance condition, and the twisted conditions. Then,*

$$\phi_{\mathcal{A}} \circ \phi_{\mathcal{S}}: \widehat{\mathcal{R}}_j \xrightarrow{\cong}^{m-1} \widehat{\mathcal{R}}_j \quad \text{and} \quad \phi_{\mathcal{S}} \circ \phi_{\mathcal{A}}: \widehat{\mathcal{R}}_{j+1} \xrightarrow{\cong}^{m-1} \widehat{\mathcal{R}}_{j+1}$$

for some $j \pmod{2}$ with $m = |j_{-1} - j_1| + 1$.

We notice that [6, Theorem 3.1] becomes a corollary of Theorem 3.3. For the proof we use the following lemma.

Lemma 3.4. *Consider*

$$K_\ell = \tilde{\phi}_{\mathcal{A}}([2\ell + 1, 2\ell + 3/2] \times [-1, 1]) \cap R_{1,0}, \quad \ell \in \mathbb{Z}$$

where $R_{1,0} = [0, 1/2] \times [-1, 1]$. If $\phi_{\mathcal{A}}$ satisfies the annular twist condition then at least $m - 1$ of the K_ℓ are non empty with $m = |j_1 - j_{-1}| + 1$.

Proof. We will prove the lemma in the case of the first annular strip condition, the proof for the second condition being similar.

Let $\theta_0 \in [0, 1/2]$ be fixed. The vertical segment (θ_0, ρ) , $\rho \in [-1, 1]$ is mapped by $\tilde{\phi}_{\mathcal{A}}$ in to a curve. Its end points satisfy

$$\begin{aligned} \tilde{\phi}_{\mathcal{A}}(\theta_0, -1) &= (\theta_{-1}, -1) \text{ where } \theta_{-1} \leq \theta_0 + 2j_{-1} + \frac{1}{2}, \\ \tilde{\phi}_{\mathcal{A}}(\theta_0, 1) &= (\theta_1, 1) \text{ where } \theta_1 \geq \theta_0 + 2j_{-1} + \frac{1}{2} + 2m - 1. \end{aligned}$$

Hence, $|\theta_{-1} - \theta_1| \geq 2m - 1$ and $K_\ell \neq \emptyset$ for $\ell = j_{-1}, \dots, j_{-1} + m - 1$. \square

Proof of Theorem 3.3. First of all without loss of generality we assume that $\phi_{\mathcal{S}}$ maps \mathcal{R}_2 across \mathcal{R}_1 thanks to the strip twist condition. Hence we prove that $\phi_{\mathcal{S}} \circ \phi_{\mathcal{A}}: \widehat{\mathcal{R}}_1 \xrightarrow{\cong}^m \widehat{\mathcal{R}}_1$. The other situations are just an adaptation of this proof.

We want to find disjoint compact subsets $\mathcal{K}_1, \dots, \mathcal{K}_{m-1} \subset \mathcal{R}_1$ such that for any continuous path γ across \mathcal{R}_1 with $\gamma(0), \gamma(1)$ in different components of $\partial\mathcal{R}_1$, the restriction $\phi_{\mathcal{A}}(\gamma(t))|_{\mathcal{K}_\ell}$ goes across \mathcal{R}_2 . In order to do this we work on the covering space, where the \mathcal{K}_ℓ will be represented by the K_ℓ of Lemma 3.4. The \mathcal{K}_ℓ are pairwise disjoint because the K_ℓ lie in a single representative $R_{1,0}$ of \mathcal{R}_1 .

The arguments used in the proof of Lemma 3.4 ensure that the curve $\tilde{\gamma}(t)$ in the covering, satisfying $\tilde{\gamma}(0) = (\theta_0, -1)$, and $\tilde{\gamma}(1) = (\theta_1, 1)$ with $\theta_0, \theta_1 \in [0, 1/2]$ goes across all the K_ℓ , and that the restriction of $\tilde{\gamma}$ to each K_ℓ goes across some copy, $[2\ell + 1, 2\ell + 3/2] \times [-1, 1]$, of R_2 . \square

4. Application to codimension 1 reversible vector fields

To detect chaotic dynamics, we apply the topological results of the previous section to some periodically forced reversible ODEs. In particular, we consider a T -periodic step-wise forcing term $p(t)$ that switches between two different values as follows

$$p(t) := \begin{cases} \lambda_1 & \text{for } t \in [0, \tau_1), \\ \lambda_2 & \text{for } t \in [\tau_1, \tau_2), \end{cases}$$

where $\lambda_1 \neq \lambda_2$ and $0 < \tau_1 < \tau_2 < T$ with $\tau_1 + \tau_2 = T$. We investigate the T -periodic problem associated with the system

$$\begin{cases} \dot{x} = yf(x, y^2), \\ \dot{y} = g(x, y^2) + p(t), \end{cases} \quad (4.1)$$

where f and g are smooth functions that identify the normal forms of codimension 1 reversible systems introduced in [11].

Our goal is to prove the existence of chaotic dynamics for the T -periodic solutions of (4.1). First, we look at the flow of the vector field $X(x, y)$ associated with (4.1) which is given by the unique solution $(x(t), y(t)) = \varphi(t, x_0, y_0)$ of $\dot{X} = X(x, y)$ satisfying $x(0) = x_0$ and $y(0) = y_0$. We study the Poincaré map $\Phi: \mathbb{R}^2 \rightarrow \mathbb{R}^2$ defined by $\Phi(x_0, y_0) = \varphi(T, x_0, y_0)$ for every point $(x_0, y_0) \in \mathbb{R}^2$. Second, we notice that the full dynamics of the problem can be broken on two sub-systems

$$\begin{cases} \dot{x} = yf(x, y^2), \\ \dot{y} = g(x, y) + \lambda_1, \end{cases} \quad (4.2)$$

and

$$\begin{cases} \dot{x} = yf(x, y^2), \\ \dot{y} = g(x, y) + \lambda_2. \end{cases} \quad (4.3)$$

Hence, we have that the Poincaré map Φ may be decomposed as $\Phi = \Phi_{\lambda_2} \circ \Phi_{\lambda_1}$, where, for any $(x_0, y_0) \in \mathbb{R}^2$, $\Phi_{\lambda_1}(x_0, y_0) = \varphi_{\lambda_1}(\tau_1, x_0, y_0)$ and $\Phi_{\lambda_2}(x_0, y_0) = \varphi_{\lambda_2}(\tau_2, x_0, y_0)$ are the Poincaré maps associated with (4.2) and (4.3), respectively. We outline here the structure of the proof for the saddle case, done by applying Theorem 3.3.

- 1) Locate a flow invariant line $\Gamma_{1,*}$ for, say λ_1 and a closed flow invariant line $\Gamma_{2,*}$ for λ_2 , making sure they intersect in at least two points. Then $\Gamma_{2,*}$ is going to be $\partial^e \mathcal{A}$ and $\Gamma_{1,*}$ will be of one component of $\partial \mathcal{S}$.
- 2) Take τ_1 to be the time it takes for φ_{λ_1} to move one intersection point to the next one.
- 3) Look at a curve γ_1 ending at the first intersection point as a candidate for a bridge and make sure Φ_{λ_1} maps it to $in(\Gamma_{2,*})$. Take P to be the other end point of γ_1 .
- 4) Take the φ_{λ_1} trajectory through P to be the other component of $\partial \mathcal{S}$ and take the (closed) φ_{λ_2} trajectory through P to be $\partial^i \mathcal{A}$. This ensures that the strip twist condition (Definition 3.5) holds.
- 5) Obtain the time τ_2 for the annular-strip condition (Definition 3.7).

In this way we can prove that the dynamics of (4.1) is semi-conjugate to a shift in a finite alphabet.

4.1. Saddle case.

We assume that system (4.1) has a saddle structure by considering

$$\begin{cases} \dot{x} = xy, \\ \dot{y} = x - y^2 + p(t). \end{cases} \tag{4.4}$$

Depending on $p(t)$, the phase-portrait of system (4.4) switches between different configurations as described in Section 2.

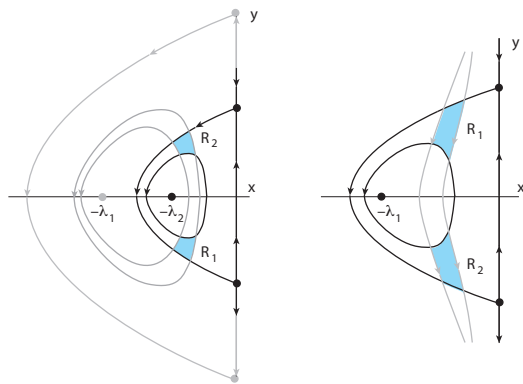


FIGURE 9. Construction of the annulus \mathcal{A} and the strip \mathcal{S} in the saddle case. Left: $0 < \lambda_2 < \lambda_1$; right: $\lambda_2 \leq 0 < \lambda_1$.

Theorem 4.1. *Let Φ be the Poincaré map associated with system (4.4). Then for each $\lambda_1 > 0$ and each λ_2 with $\lambda_1 > \lambda_2$ and for an open set of values of τ_1 and τ_2 the map Φ induces chaotic dynamics on $m \geq 2$ symbols.*

Proof. First of all we notice that the following two cases can occur: $\lambda_1 > \lambda_2 > 0$ or $\lambda_1 > 0 \geq \lambda_2$.

Let us suppose that λ_1 and λ_2 are two fixed positive values satisfying the first case. Then for both systems (4.2) and (4.3) there exist three equilibria. In particular, there exists a heteroclinic cycle around the center $(-\lambda_i, 0)$ which joins the two saddles $(0, -\sqrt{\lambda_i})$ and $(0, \sqrt{\lambda_i})$, for $i = 1, 2$.

Let $(x^*, 0)$ be the point where the heteroclinic cycle of system (4.3) crosses the negative part of the x -axis. Then two configurations are possible: $-\lambda_1 < x^* < -\lambda_2$ or $x^* < -\lambda_1$. It will be not restrictive to consider the first configuration since the other situation can be treated similarly. We proceed with the construction of an annulus \mathcal{A} and a strip \mathcal{S} which satisfy the topological conditions required to apply Theorem 3.3.

For any $(x, y) \in \mathbb{R}^2$, we call $\Gamma_1(x, y)$ and $\Gamma_2(x, y)$ the trajectories through the point (x, y) of system (4.2) and (4.3), respectively. Let $\Gamma_2(x^*, 0)$ be the heteroclinic trajectory through $(x^*, 0)$, then we define the outer component of

$\partial\mathcal{A}$ as $\partial^e\mathcal{A} := \Gamma_2(x^*, 0) \cup \{(0, -\sqrt{\lambda_2})\} \cup \Gamma_2(0, 0) \cup \{(0, \sqrt{\lambda_2})\}$. Let $\alpha < 0$ with $-\lambda_2 < \alpha$ be any number so the trajectory $\Gamma_1(\alpha, 0)$ through $(\alpha, 0)$ will cross the heteroclinic connection $\Gamma_2(x^*, 0)$. We take $\Gamma_1(\alpha, 0) \cap \{x^* \leq x \leq 0\}$ to be one of the components of $\partial\mathcal{S}$, and we construct the other two boundary pieces of the annulus and the strip so as to satisfy the linkage condition and the twist conditions.

Let τ_1 be the minimum positive time such that, if $p(t)$ is a solution of (4.2) through $(\alpha, 0)$ with $p(0) \in \Gamma_2(x^*, 0) \cap \{y < 0\}$, then $p(\tau_1) \in \Gamma_2(x^*, 0) \cap \{y > 0\}$. For any point $(x, y) \in \Gamma_2(x^*, 0) \cap \{y < 0\}$ close to $p(0)$ the points $\varphi_{\lambda_1}(\tau_1, x, y)$ form a curve through $p(\tau_1)$. Generically this curve goes across $\Gamma_2(x^*, 0)$ (otherwise, make a small change in α). Suppose that the curve is below $\Gamma_2(x^*, 0)$ to the left of $p(\tau_1)$ (otherwise the arguments are similar). Take $\beta < 0$ with $-\lambda_2 < \beta < \alpha < 0$ such that the points in the trajectory $\Gamma_1(\beta, 0)$ of system (4.2) through $(\beta, 0)$ satisfy the condition on the curve. Then we take the other component of $\partial\mathcal{S}$ as $\Gamma_1(\beta, 0) \cap \{x^* \leq x \leq 0\}$. It remains to obtain the inner component of $\partial\mathcal{A}$.

Let $\Pi: \mathbb{R}^2 \rightarrow \mathbb{R}^2$ be the projection on the second component, namely $\Pi(x, y) = y$. For any $(x, y) \in \mathbb{R}^2$ let $\psi(x, y) = \Pi(\varphi_{\lambda_1}(\tau_1, x, y))$ and let $\bar{\psi}(x, y) = \psi(x, y) + \Pi(x, y)$, so $\bar{\psi}(x, y)$ compares the height of $\varphi_{\lambda_1}(\tau_1, x, y)$ to that of the symmetric point of (x, y) .

Let $q(t)$ be the solution of (4.2) through $(\beta, 0)$ with $q(0) \in \Gamma_2(x^*, 0) \cap \{y < 0\}$. Then $\bar{\psi}(q(0)) < 0$. Also there exists a $\sigma > 0$ such that $q(\sigma) \in \Gamma_2(x^*, 0) \cap \{y > 0\}$. By construction, $\bar{\psi}(q(\sigma)) > 0$. Therefore, there exists $\hat{\sigma} \in (0, \sigma)$ such that $\bar{\psi}(q(\hat{\sigma})) = 0$. This means that $\varphi_{\lambda_1}(\tau_1, q(\hat{\sigma}))$ is symmetric to $q(\hat{\sigma})$. The trajectory $\Gamma_2(q(\hat{\sigma}))$ will go through both $q(\hat{\sigma})$ and $\varphi_{\lambda_1}(\tau_1, q(\hat{\sigma}))$. We define the inner component of $\partial\mathcal{A}$ as $\partial^i\mathcal{A} := \Gamma_2(q(\hat{\sigma}))$.

In this manner, the topological annulus \mathcal{A} and the topological strip \mathcal{S} are linked by construction (see Figure 9). The linkage condition gives two symmetric topological rectangles \mathcal{R}_1 and \mathcal{R}_2 (in the lower and upper half-plane, respectively) that satisfy the twist conditions. Indeed, a strip-twist condition holds for $\Phi_{\lambda_1}: \mathcal{S} \rightarrow \mathcal{S}$ because the rectangle $\mathcal{R}_1 \subset \mathcal{A} \cap \mathcal{S} \cap \{y < 0\}$ is stretched across $\mathcal{R}_2 \subset \mathcal{A} \cap \mathcal{S} \cap \{y > 0\}$. Since $\Gamma_2(x^*, 0)$ is a heteroclinic connection then for τ_2 large enough also an annulus-twist condition holds for $\Phi_{\lambda_2}: \mathcal{A} \rightarrow \mathcal{A}$ because \mathcal{R}_2 is stretched across \mathcal{R}_1 a suitable number of times (depending on τ_2). The result follows by an application of Theorem 3.3 to the Poincaré map $\Phi = \Phi_{\lambda_2} \circ \Phi_{\lambda_1}$. This concludes the first case.

The proof above holds for a fixed value of τ_1 and for sufficiently large τ_2 . However, we may obtain the result for τ_1 in an open interval by taking different value of α .

The arguments above yield a proof for the case $\lambda_1 > 0 \geq \lambda_2$, we just indicate where it needs to be adapted. The outer component of $\partial\mathcal{A}$ may be taken as $\partial^e\mathcal{A} := \Gamma_1(x^*, 0) \cup \{(0, -\sqrt{\lambda_1})\} \cup \Gamma_1(0, 0) \cup \{(0, \sqrt{\lambda_1})\}$, where $\Gamma_1(x^*, 0)$ is the heteroclinic trajectory of φ_{λ_1} going through $(x^*, 0)$. One of the components of $\partial\mathcal{S}$ will be $\Gamma_2(\alpha, 0)$ with $-\lambda_1 < \alpha < 0$.

Then take τ_2 to be the least positive time to go from $\Gamma_2(\alpha, 0) \cap \Gamma_1(x^*, 0) \cap \{y > 0\}$ to $\Gamma_1(x^*, 0) \cap \{y < 0\}$. Apply the arguments above to obtain the other component of $\partial\mathcal{S}$ as a φ_{λ_2} trajectory that starting at $\Gamma_1(x^*, 0) \cap \{y > 0\}$ arrives above $\Gamma_1(x^*, 0) \cap \{y < 0\}$ in time τ_2 . Then find a point q in this trajectory and in the upper half-plane, such that Φ_{λ_2} maps q to its symmetric $h(q)$. Take $\partial^i\mathcal{A} := \Gamma_2(q)$ to complete the construction. \square

4.2. Cusp case.

When system (4.1) has the following form

$$\begin{cases} \dot{x} = y, \\ \dot{y} = x^2 + p(t). \end{cases} \quad (4.5)$$

then its phase-portrait is of cusp type. We notice that system (4.5) has also a Hamiltonian structure, and in this juncture, when $\lambda_1 < 0$ and $\lambda_2 \leq 0$ the geometry is similar to the one investigated in [9, 10]. Hence, we expect that chaotic dynamics occurs for τ_1 and τ_2 large enough. On the other hand, by applying the procedure exploited for Theorem 4.1, we can prove what follows.

Theorem 4.2. *Let Φ be the Poincaré map associated with system (4.5). Then for each $\lambda_1 \leq 0$ and each λ_2 with $\lambda_1 < \lambda_2$ and for an open set of values of τ_1 and τ_2 the map Φ induces chaotic dynamics on $m \geq 2$ symbols.*

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Isabel S. Labouriau

Centro de Matemática da Universidade do Porto
Rua do Campo Alegre 687, 4169-007 Porto, Portugal
e-mail: islabour@fc.up.pt

Elisa Sovrano

Centro de Matemática da Universidade do Porto
Rua do Campo Alegre 687, 4169-007 Porto, Portugal
e-mail: elisasovrano@gmail.com